

Anisotropic Poroelasticity in Cracked and Granular Materials

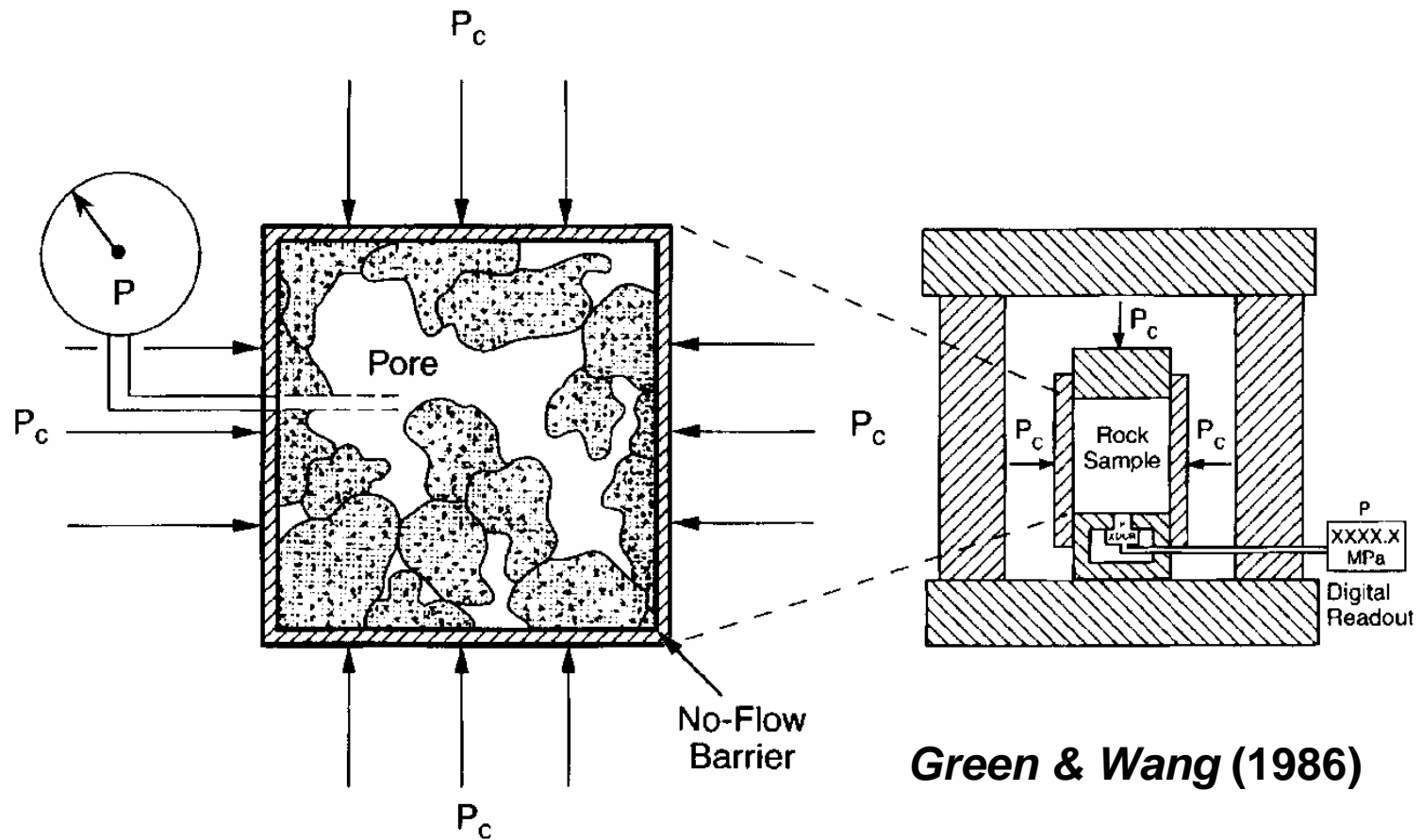


Teng-fong Wong
Nicholas M. Beeler
David A. Lockner

Anisotropic Poroelasticity in Cracked and Granular Materials

- Undrained deformation, the Skempton coefficient and its use in the analysis of earthquakes
- Laboratory measurement of anisotropic poroelastic coefficients in crustal rocks
- Effects of anisotropic cracking on the Skempton tensor





Green & Wang (1986)

Undrained Condition $\rightarrow \Delta m = 0$

Skempton's Coefficient: $B = \left. \frac{dp}{dP_c} \right|_{\Delta m=0}$

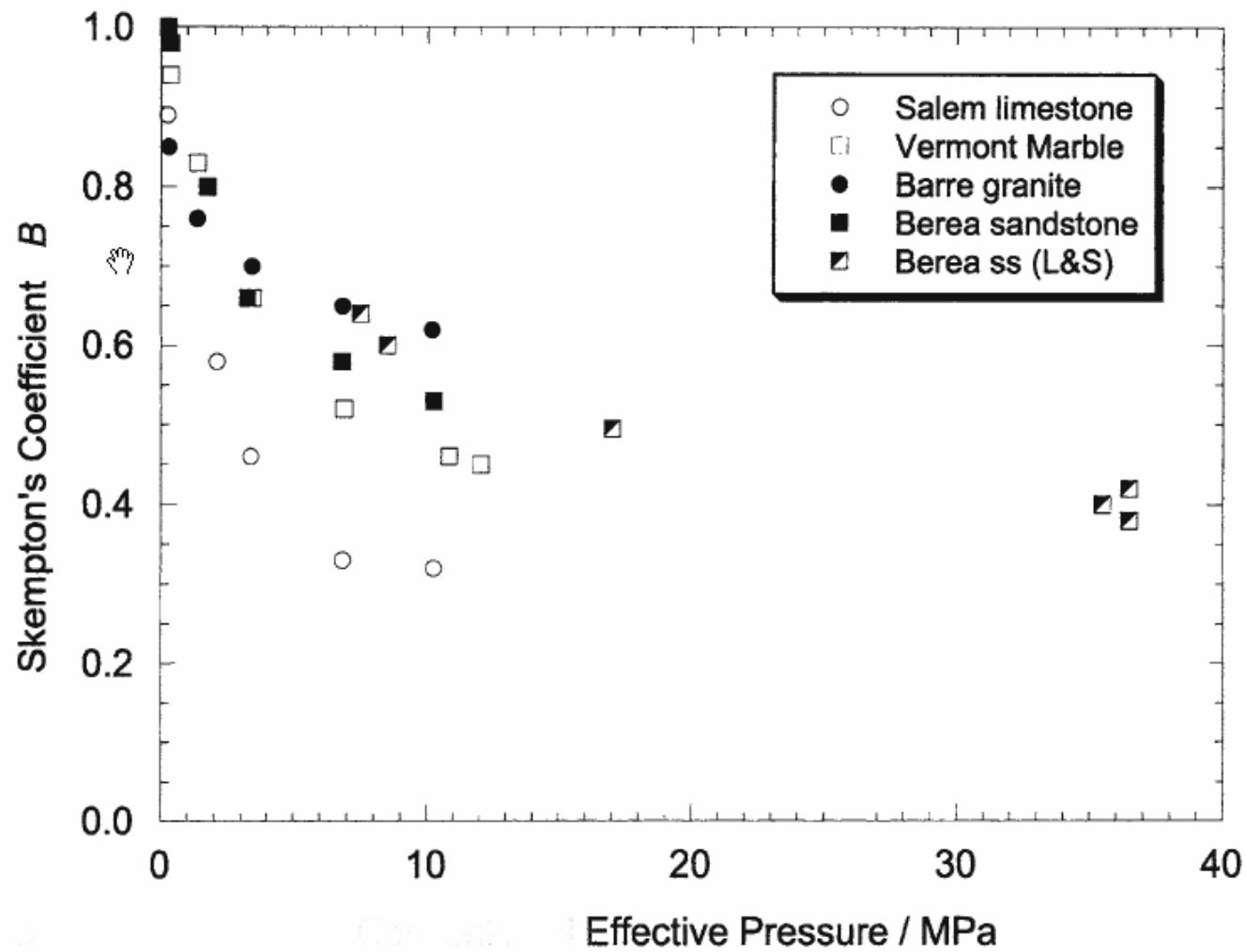
TABLE C.1.

Example Poroeleastic Moduli for Different Rock Types*

Rock	G (GPa)	ν	K (GPa)	K'_s (GPa)	α	ν_u	K_u (GPa)	B	c (m ² /s)	ϕ	k (mD)
Berea sandstone	6	0.20	8.0	36	0.79	0.33	16.0	0.62	1.6×10^0	0.19	1.9×10^2
Boise sandstone	4.2	0.15	4.6	42	0.85	0.31	8.3	0.50	4.0×10^{-1}	0.26	8.0×10^2
Ohio sandstone	6.8	0.18	8.4	31	0.74	0.28	13.0	0.50	3.9×10^{-2}	0.19	5.6×10^0
Pecos sandstone	5.9	0.16	6.7	39	0.83	0.31	14.0	0.61	5.4×10^{-3}	0.20	8.0×10^{-1}
Ruhr sandstone	13	0.12	13	36	0.65	0.31	30.0	0.88	5.3×10^{-3}	0.02	2.0×10^{-1}
Weber sandstone	12	0.15	13	36	0.64	0.29	25.0	0.73	2.1×10^{-2}	0.06	1.0×10^0
Tennessee marble	24	0.25	40	50	0.19	0.27	44.0	0.51	1.3×10^{-5}	0.02	1.0×10^{-4}
Charcoal granite	19	0.27	35	45	0.27	0.30	41.0	0.55	7.0×10^{-6}	0.02	1.0×10^{-4}
Westerly granite	15	0.25	25	45	0.47	0.34	42.0	0.85	2.2×10^{-5}	0.01	4.0×10^{-4}
Clay	—	—	0.062	∞	1	—	6.2	0.99	—	—	—
Mudstone	—	—	2.13	42	0.95	—	10.1	0.83	—	—	—
Kayenta sandstone	—	—	9.1	37.9	0.76	—	18.5	0.67	—	—	—
Limestone	—	—	33.3	107	0.69	—	40.2	0.25	—	—	—
Hanford basalt	—	—	46.7	59	0.23	—	45.4	0.12	—	—	—
Berea sandstone	5.6	0.17	6.6	28.9	0.77	0.34	15.8	0.75	1.5×10^0	0.19	1.9×10^2
Indiana limestone	12.1	0.26	21.2	72.6	0.71	0.33	31.2	0.46	—	0.13	—

*The top section of the table is a combination of measured and calculated values compiled by Detournay and Cheng (1993), who incorporated values previously presented by Rice and Cleary (1976). Typically, drained moduli are measured values, and the solid-grain modulus is taken from handbooks of mineral properties. The undrained moduli are computed assuming that the solid-grain modulus K_s is equal to both theunjacketed bulk modulus K'_s and theunjacketed pore incompressibility K_ϕ . The fluid-bulk modulus was assumed to be $K_f = 3.3$ GPa. The second section of the table compiled by Palciauskas and Domenico (1989) and reprinted in Domenico and Schwartz (1998, p. 171) also combined measurement and calculation. The bottom set of values for Berea sandstone and Indiana limestone is a best-fit set of constants based on eight different measurements by Hart and Wang (1995) in which water ($K_f = 2.3$ GPa) was the pore fluid.

Wang (2000)



Paterson & Wong (2005)

Constitutive Relations for an Isotropic Poroelastic Material (Biot, 1941; Rice & Cleary, 1976)

$$2G\varepsilon_{ij} = \sigma_{ij} - \frac{\nu}{1 + \nu} \sigma_{kk} \delta_{ij} + \frac{3(\nu_u - \nu)}{B(1 + \nu)(1 + \nu_u)} p \delta_{ij}$$

$$m = m_0 + \frac{3\rho_0(\nu_u - \nu)}{2GB(1 + \nu)(1 + \nu_u)} \left(\sigma_{kk} + \frac{3}{B} p \right)$$

Undrained Condition ($\Delta m = 0$ or $m = m_0$) \rightarrow

$$\Delta p = -B \Delta \sigma_{kk} / 3$$

$$2G \Delta \varepsilon_{ij} = \Delta \sigma_{ij} - \frac{\nu_u}{1 + \nu_u} \Delta \sigma_{kk} \delta_{ij}$$

(Tensile stress is considered to be positive here)

➤ In an **isotropic** poroelastic material, the **pore pressure induced by undrained deformation** is proportional to the **mean stress** and solely characterized by the scalar **Skempton coefficient**:

$$\Delta p = B \frac{\Delta \sigma_{kk}}{3} = B \Delta \sigma_m$$

➤ The induced pore pressure is **independent of the deviatoric stresses** in an isotropic poroelastic material.

(Compressive stress is considered to be positive from this point on)

Coulomb Stress Analysis of Earthquake Interaction: Poroelastic effects due to undrained pore pressure excess (*Beeler et al., 2000*)

$$CFF = \tau - \mu(\sigma_n - p)$$

$$\Delta CFF = \Delta\tau - \mu(\Delta\sigma_n - \Delta p)$$

The deformation is undrained, so the induced pore pressure is characterized by the Skempton coefficient

$$\Delta CFF = \Delta\tau - \mu(\Delta\sigma_n - B\Delta\sigma_m)$$

$$= \Delta\tau - \mu \left(1 - B \frac{\Delta\sigma_m}{\Delta\sigma_n}\right) \Delta\sigma_n = \Delta\tau - \mu' \Delta\sigma_n$$

$$\text{with } \mu' = \mu \left(1 - B \frac{\Delta\sigma_m}{\Delta\sigma_n}\right)$$

The mean stress distribution in the vicinity of a dislocation source is highly heterogeneous, and yet in simulations it is often assumed that μ' is constant: $\mu' \approx \mu (1 - B)$?

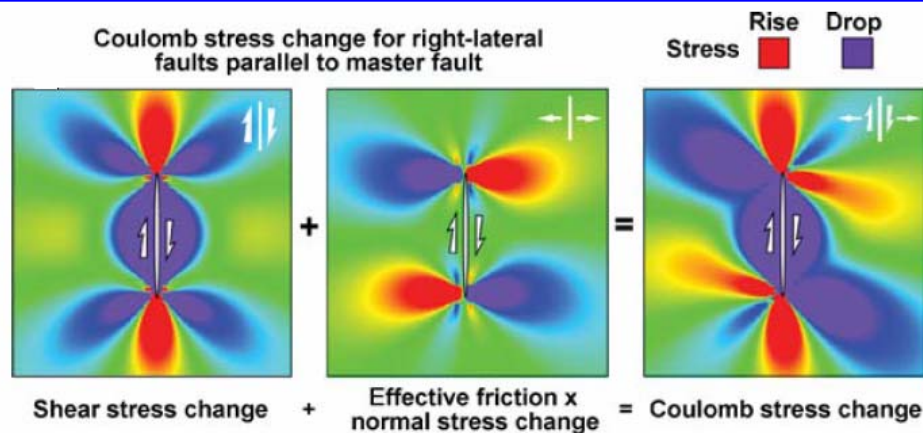
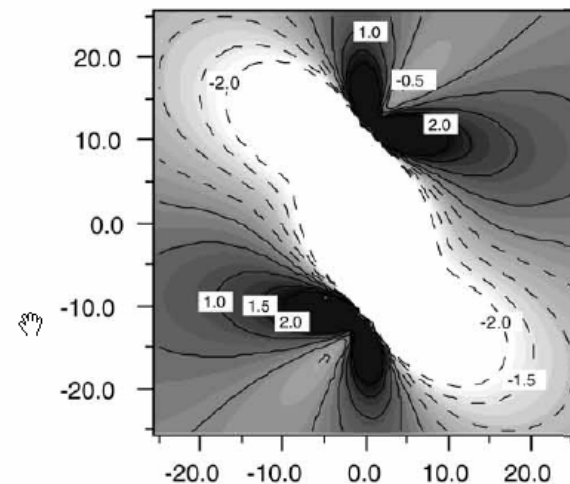
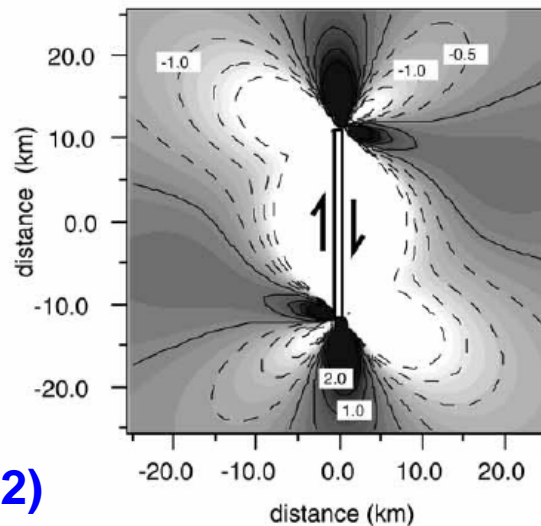


Figure 2 Illustration of a Coulomb stress change calculation (Equation 4). The panels show a map view of a vertical strike-slip fault embedded in an elastic half-space, with imposed slip that tapers toward the fault ends. Stress changes are depicted by graded colors; green represents no change in stress. (From King et al. 1994.)

constant apparent friction
 $\Delta CFF = \Delta\tau + \mu'\Delta\sigma_n$

$\Delta CFF = \Delta\tau - \mu(\Delta\sigma_n - B\Delta\sigma_m)$

$\mu' = 0.4; \mu = 0.75; B = 0.47$

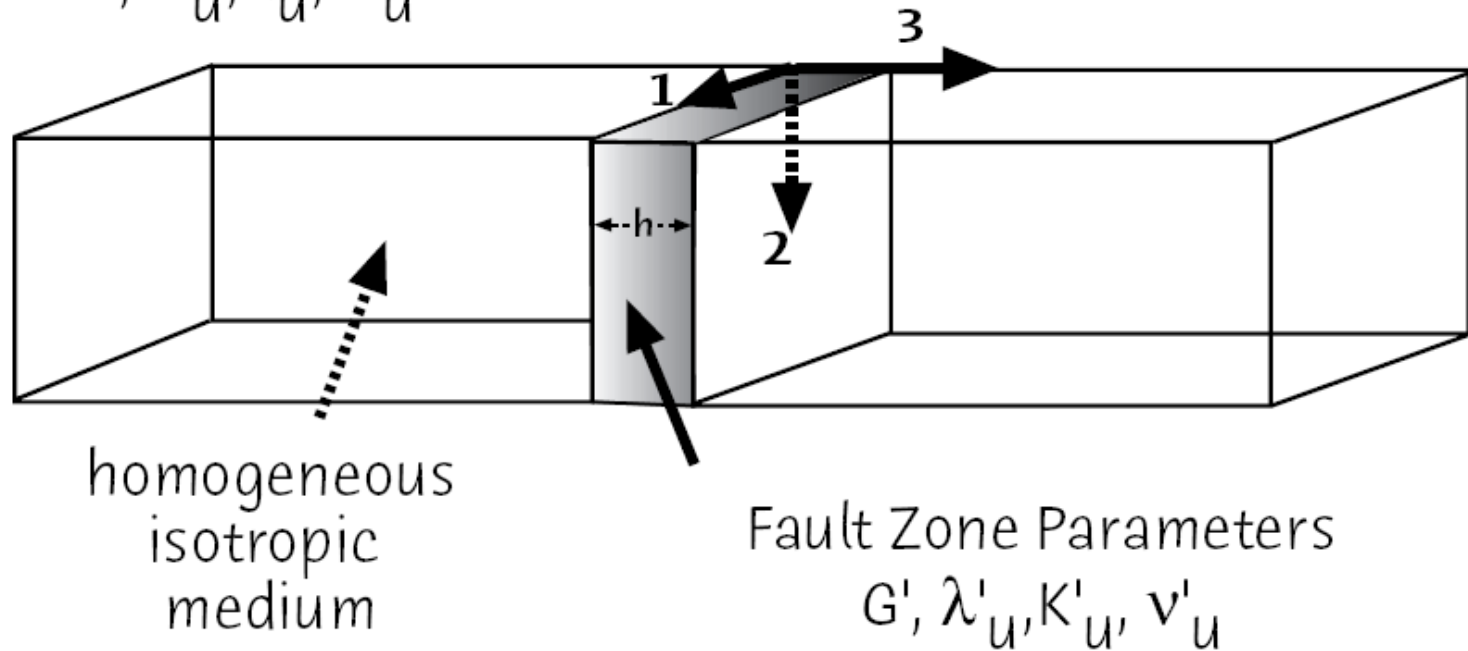


Cocco & Rice (2002)

Surrounding Crust

G, λ_u, K_u, ν_u

Cocco & Rice (2002)



The fault zone and country rock have different poroelastic properties. Mechanical equilibrium and strain continuity \rightarrow pore pressure excess induced under undrained condition is given by

$$\Delta p = B' \frac{(1 + \nu'_u)}{(1 - \nu'_u)} \left[\frac{G'}{G} \frac{(1 - \nu_u)}{(1 + \nu_u)} \frac{\Delta \sigma_{kk}}{3} + \frac{(G - G')}{G} \Delta \sigma_{33} \right]$$

Pore pressure excess that develops under undrained condition in a fault zone with poroelastic properties different from the surrounding.

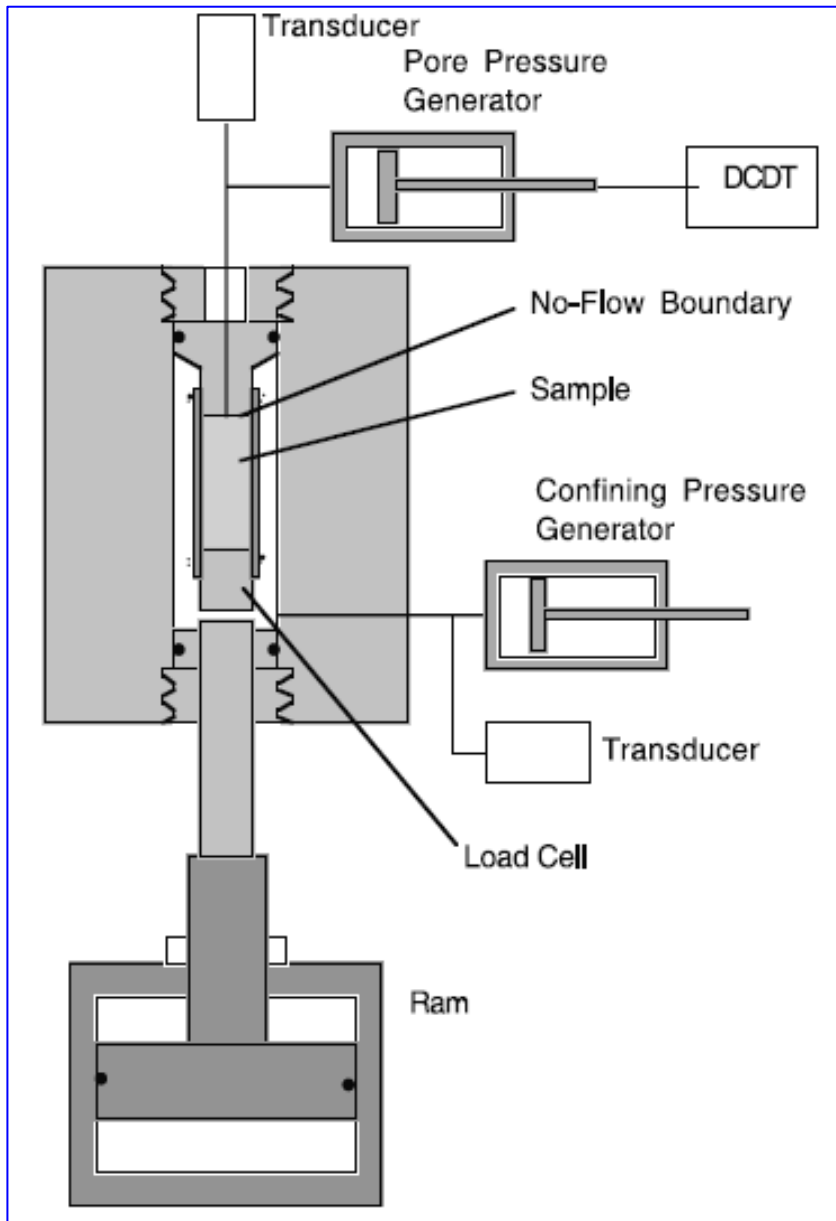
$$\Delta p = B' \frac{(1 + \nu'_u)}{(1 - \nu'_u)} \left[\frac{G'}{G} \frac{(1 - \nu_u)}{(1 + \nu_u)} \frac{\Delta \sigma_{kk}}{3} + \frac{(G - G')}{G} \Delta \sigma_{33} \right]$$

Three limiting cases:

(1) $G' = G$
$$\Delta p = B' \frac{(1 + \nu'_u)}{(1 - \nu'_u)} \frac{(1 - \nu_u)}{(1 + \nu_u)} \frac{\Delta \sigma_{kk}}{3}$$

(2) $G \gg G'$
$$\Delta p = B' \frac{(1 + \nu'_u)}{(1 - \nu'_u)} \Delta \sigma_{33}$$

(3) The undrained response is **anisotropic**, characterized by a **Skempton tensor** B_{ij} . If the tensor component B_{33} is significantly larger than the others, then the pore pressure change is dominated by $(B_{33}/3) \Delta \sigma_{33}$.

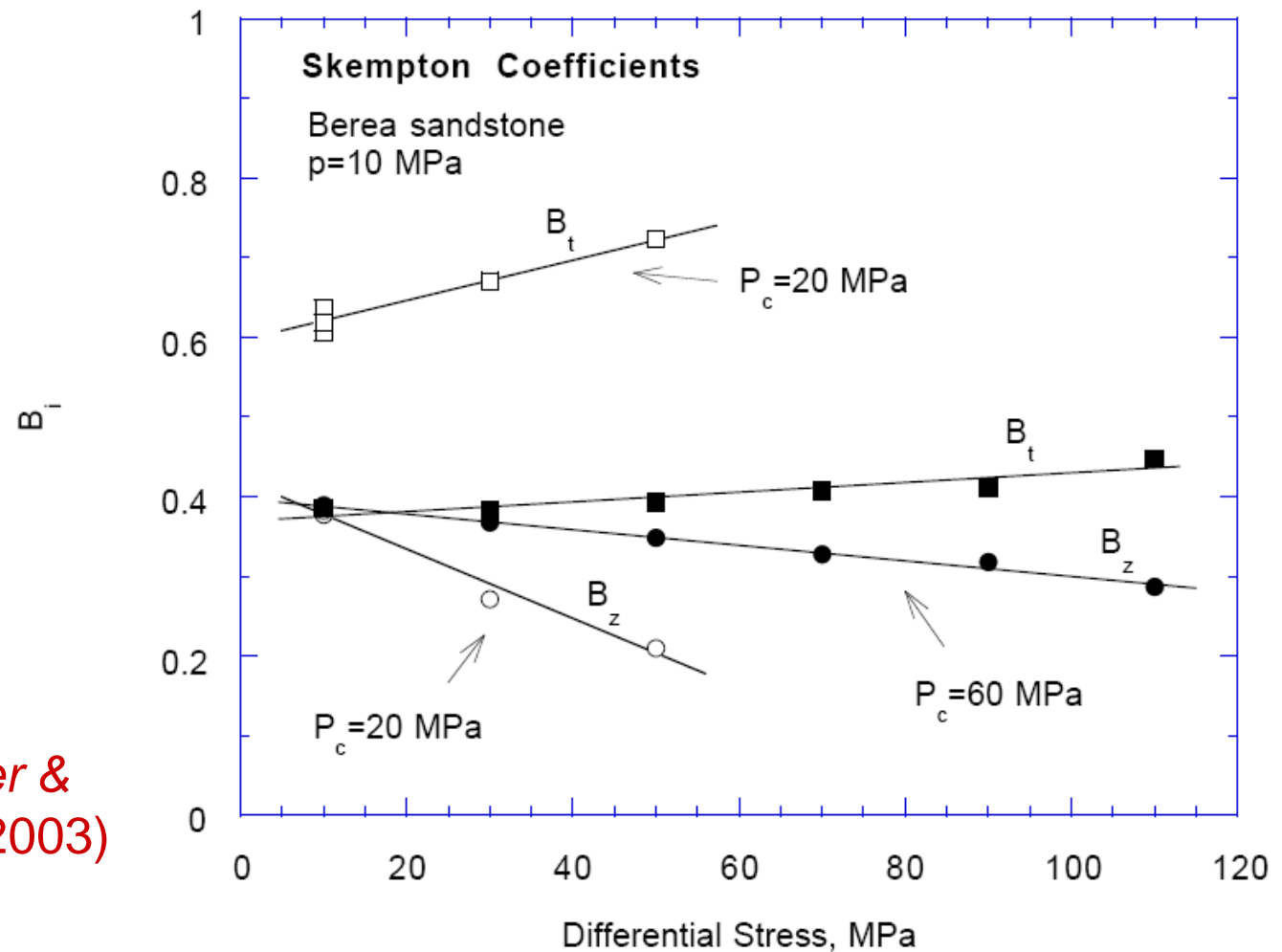


The **Skempton tensor** is defined by:

$$\Delta p = B_{ij} \frac{\Delta \sigma_{ij}}{3} \bigg|_{\Delta m=0}$$

- ❖ Measurement of pore pressure change induced by perturbations in the principal stresses.
- ❖ Storage in the pore pressure system changes in response to pore pressure, which has to be accounted for by a servo-controlled pore pressure generator.
- ❖ Systematic measurement of the Skempton tensor in triaxially compressed samples of Berea sandstone, Navajo sandstone and Ottawa sand.

Lockner and Stanchits (2002)



*Lockner &
Beeler (2003)*

- ❖ The poroelastic anisotropy is such that the undrained response to the **transverse** stress is more **pronounced**.
- ❖ The anisotropy **increases** with increasing **differential stress** and **decreases** with increasing **confinement**.

Anisotropic Poroelasticity

(Biot, 1955; Carroll, 1979; Thompson and Willis, 1991; Cheng, 1997)

$$B_1 = B_2 = \frac{[3(S_{11} + S_{12} + S_{13}) - \beta_s]}{C}$$
$$B_3 = \frac{[2S_{13} + S_{33}) - \beta_s]}{C}$$

- ❖ Transversely isotropic material with symmetry axis along the x_3 direction.
- ❖ The “drained” effective compliance of the anisotropic poroelastic material is given by S_{ij} (in the Voigt matrix notation)
- ❖ Compressibility of the solid grains is denoted by β_s (Solid assumed to be isotropic).
- ❖ The storage coefficient C is defined to be the volume of fluid released from unit rock volume per unit decrease in pore pressure under the condition of constant confining stresses.

Compliance of an Anisotropically Cracked Rock

(Kachanov, 1993; Sayers and Kachanov, 1995)

$$S_{11} = S_{22} = \frac{1}{E_o} + \alpha_{11} + \beta_{1111}$$

$$S_{33} = \frac{1}{E_o} + \alpha_{33} + \beta_{3333}$$

$$S_{12} = -\frac{\nu_o}{E_o} + \frac{\beta_{1111}}{3}$$

$$S_{23} = S_{13} = -\frac{\nu_o}{E_o} + \beta_{1133}$$

- ❖ Transversely isotropic material with an axisymmetric distribution of microcracks.
- ❖ Solid matrix is assumed to be isotropic, with elastic moduli ν_o and E_o .
- ❖ Crack fabric is characterized by two tensors: α_{ij} and β_{ijkl} .

Skempton Tensor of an Anisotropically Cracked Rock

$$B_1 = B_2 = \frac{32}{E_o C} \frac{(1 - \nu_o^2)}{(2 - \nu_o)} \left(\alpha_{11}^* + \cancel{\frac{4}{3} \beta_{1111}^*} + \cancel{\beta_{1133}^*} \right)$$

$$\Rightarrow B_t \approx \frac{32(1 - \nu_o^2)}{E_o C (2 - \nu_o)} \alpha_{11}^*$$

$$B_3 = \frac{32(1 - \nu_o^2)}{E_o C (2 - \nu_o)} \left(\alpha_{33}^* + \cancel{\beta_{3333}^*} + \cancel{2\beta_{1133}^*} \right)$$

$$\Rightarrow B_z \approx \frac{32(1 - \nu_o^2)}{E_o C (2 - \nu_o)} \alpha_{33}^*$$

- ❖ The fourth-rank crack fabric tensor can be neglected in a saturated rock under drained conditions for $\nu_o \ll 2$.
- ❖ The Skempton tensor is directly proportional to the second-rank crack fabric tensor.

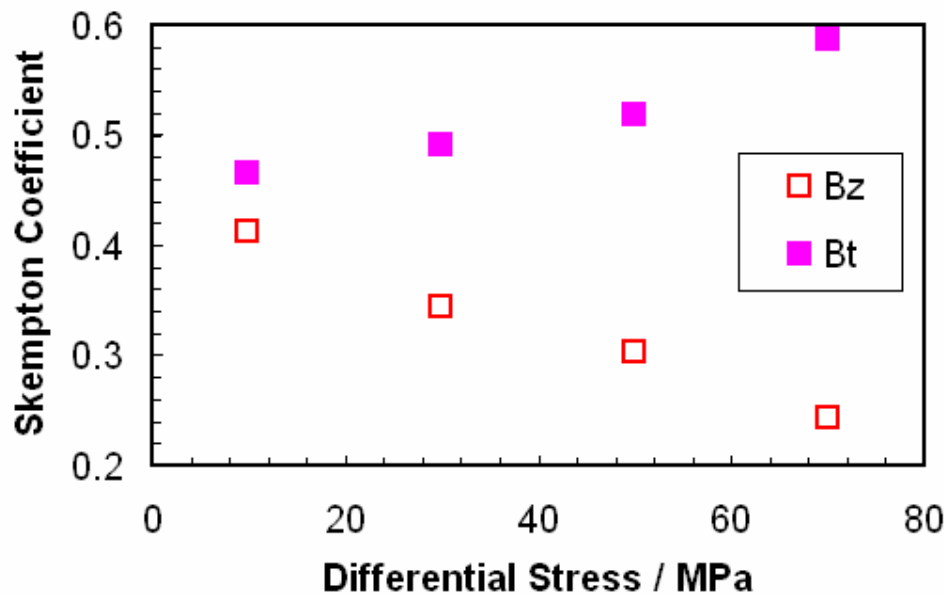
Crack Fabric Tensor and Crack Density

$$\alpha_{ij} = \frac{1}{V} \sum_{r=1}^N B_T^r n_i^r n_j^r S_r = \frac{32}{3E_o} \frac{(1 - \nu_o^2)}{(2 - \nu_o)} \alpha_{ij}^*$$

$$\alpha_{ij}^* = \frac{c^3}{V} \sum_{r=1}^N n_i^r n_j^r = \alpha_{mm}^* \langle n_i n_j \rangle$$

$$\alpha_{mm}^* = \frac{N c^3}{V}$$

- ❖ The **trace of the second-rank crack fabric tensor** is identical to the **scalar crack density** introduced by *Walsh* (1965) and *Budiansky and O'Connell* (1976).
- ❖ Indices of the tensor are associated with directions of the **crack normal** and **displacement discontinuity**.

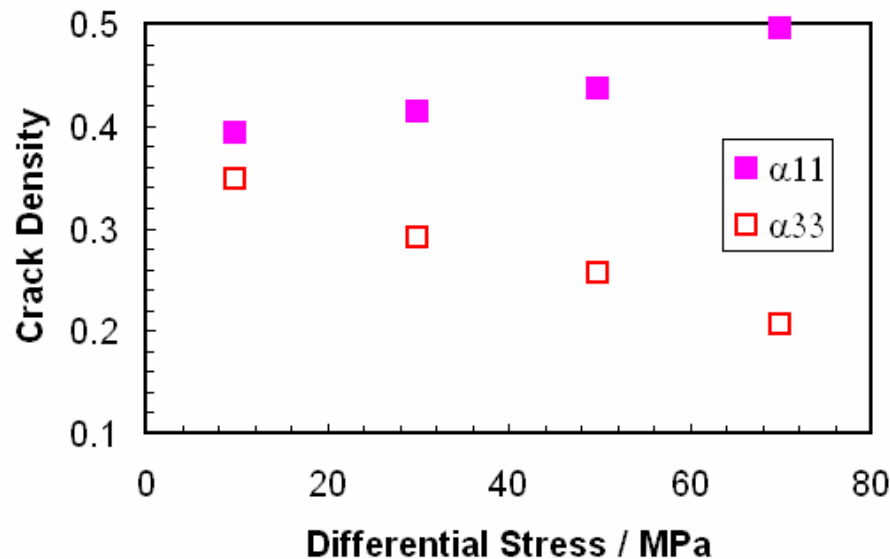


Berea sandstone data of
Lockner & Beeler (2003)

Confining Pressure: 30 MPa
Pore Pressure: 7.5 MPa

$E_o = 75$ GPa, $\nu_o = 0.15$

$C = 0.19$ GPa⁻¹ (from direct
storage measurement)



$$B_t = \frac{32(1 - \nu_o^2)}{E_o C (2 - \nu_o)} \alpha_{11}^*$$

$$B_z = \frac{32(1 - \nu_o^2)}{E_o C (2 - \nu_o)} \alpha_{33}^*$$

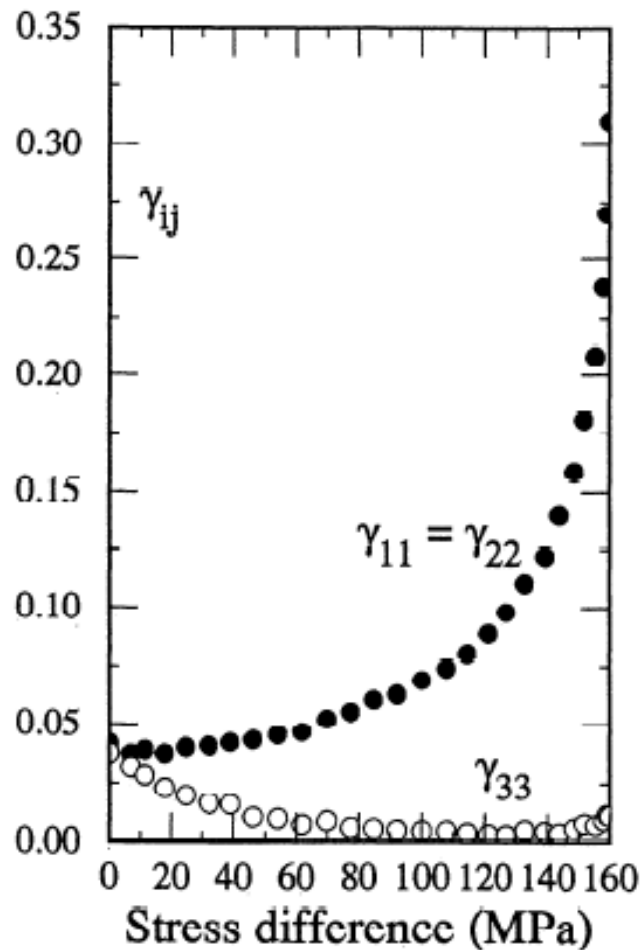


Figure 8. Components $\gamma_{11} = \gamma_{22}$ (solid dots) and γ_{33} (open dots) of the crack density tensor obtained by inverting the ultrasonic measurements of *Scott et al.* [1993] on a cylinder of Berea sandstone at 20 MPa confining pressure. Ox_3 is parallel to the maximum compressive stress.

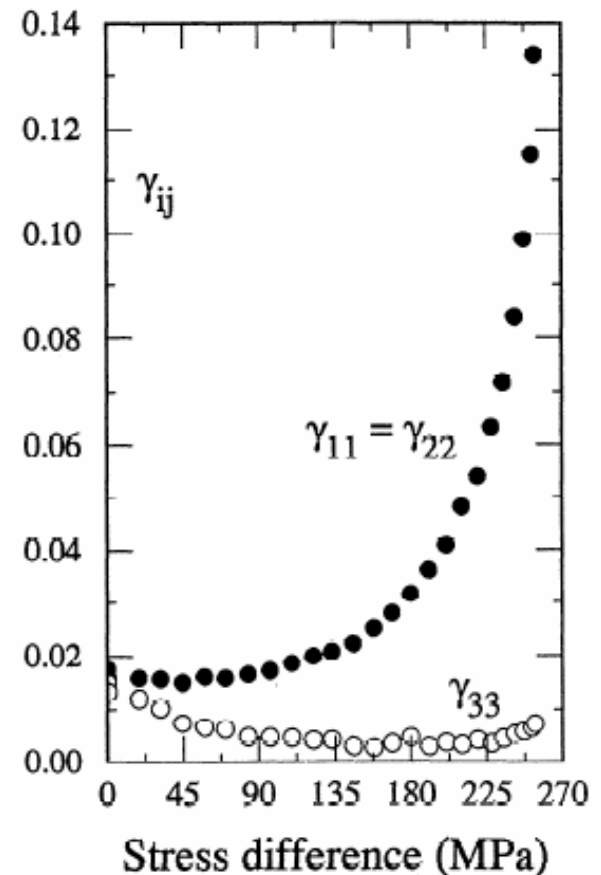
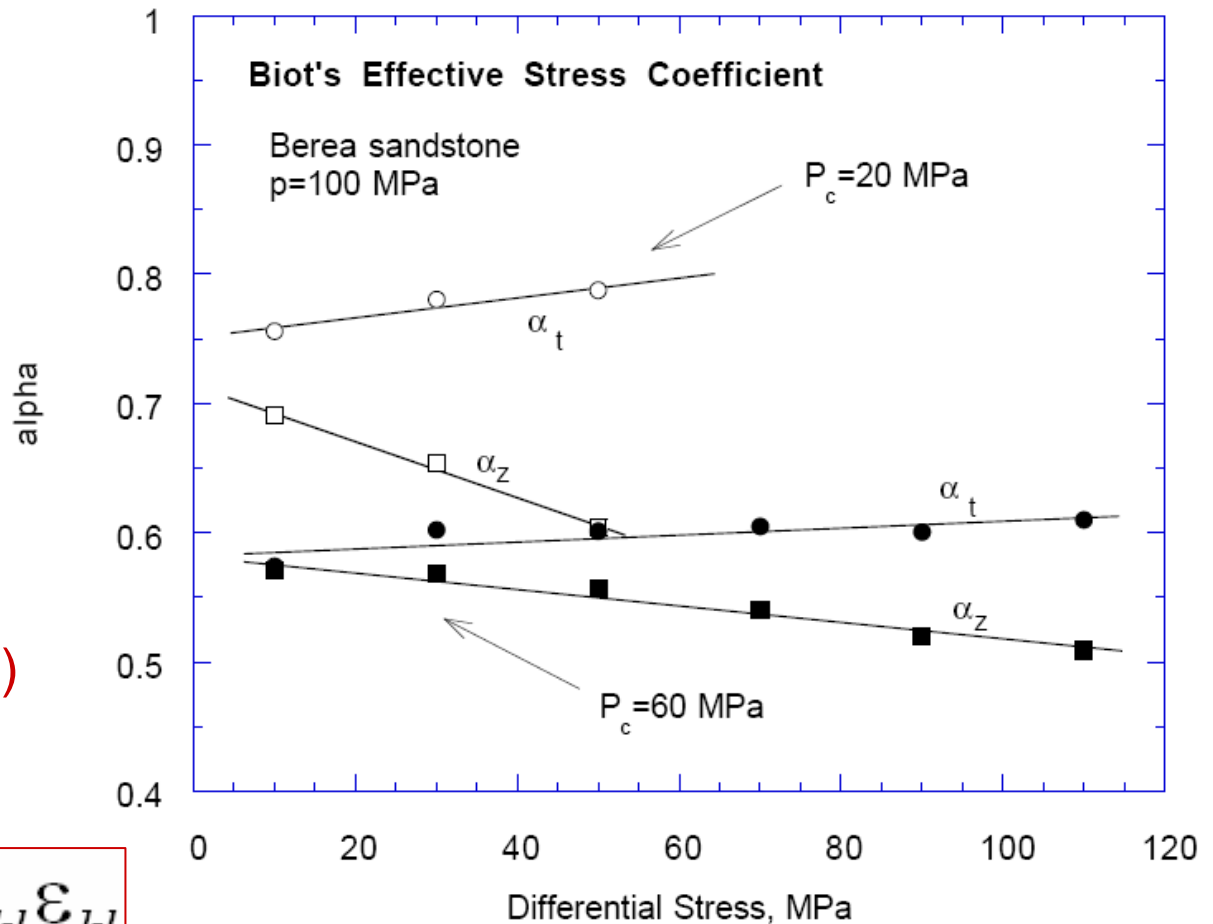


Figure 9. Components $\gamma_{11} = \gamma_{22}$ (solid dots) and γ_{33} (open dots) of the crack density tensor obtained by inverting the ultrasonic measurements of *Scott et al.* [1993] on a cylinder of Berea sandstone at 60 MPa confining pressure as a function of the difference between the maximum compressive stress and the confining pressure. Ox_3 is parallel to the maximum compressive stress.

*Lockner &
Beeler (2003)*



$$\sigma_{ij} - \alpha_{ij} p = C_{ijkl} \varepsilon_{kl}$$

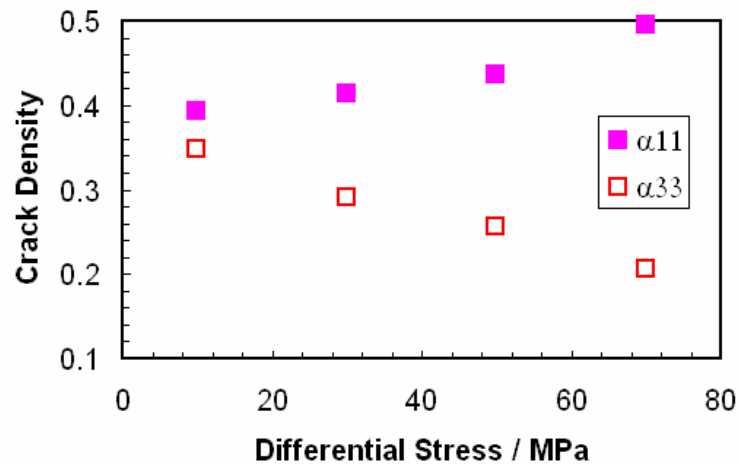
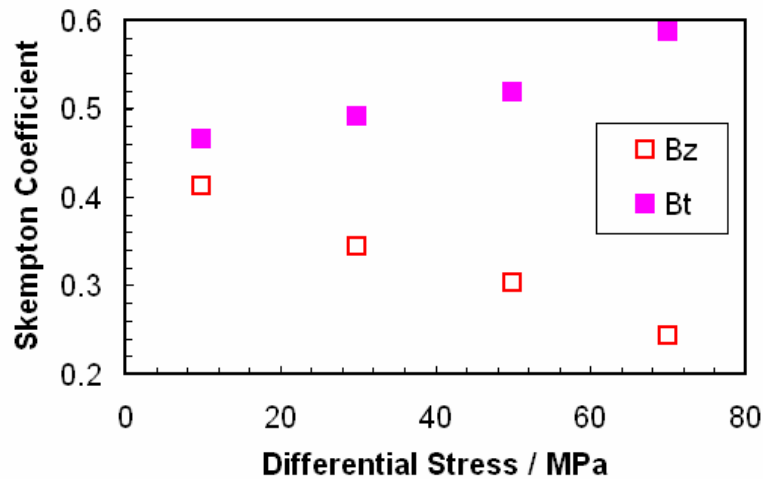
- ❖ The poroelastic anisotropy is such that the Biot-Willis coefficient is **larger** in response to the **transverse** stress.
- ❖ The anisotropy **increases** with increasing **differential stress** and **decreases** with increasing **confinement**.

Biot-Willis Effective Stress Coefficients of an Anisotropic Cracked Rock

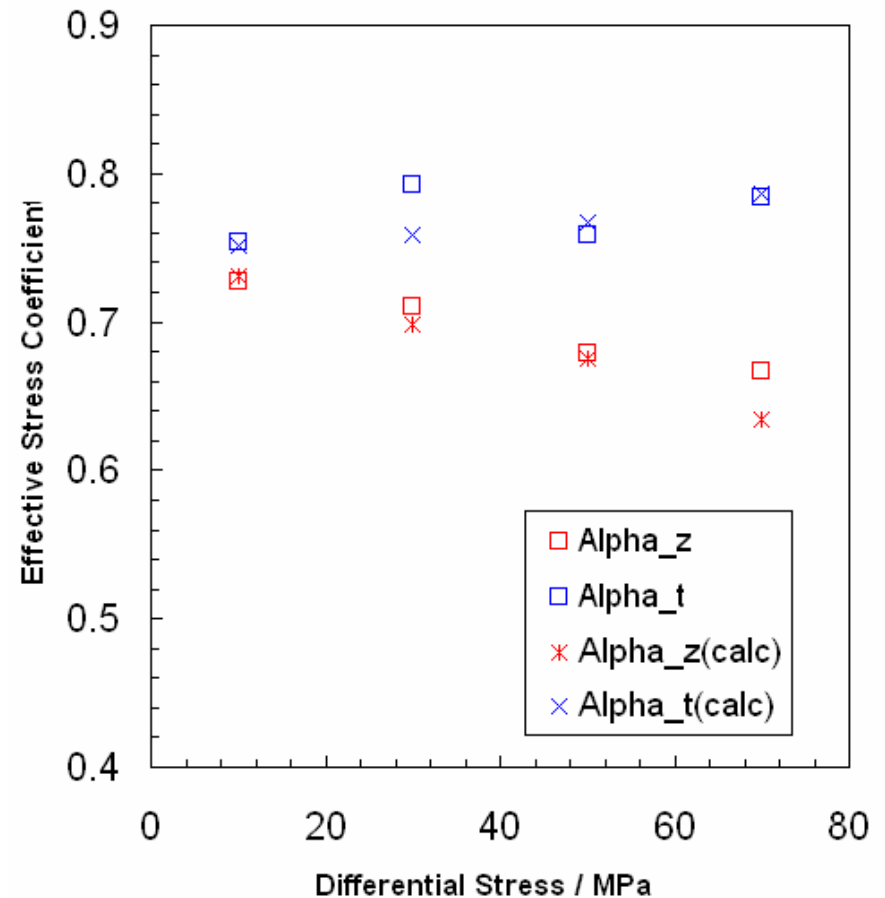
$$\alpha_t = \alpha_1 = \alpha_2 = 1 - \frac{(1 - 2\nu_o)(1 + \nu_o + g\alpha_{33}^*)}{(1 + g\alpha_{33}^*)(1 - \nu_o + g\alpha_{11}^*) - 2\nu_o^2}$$

$$\alpha_z = \alpha_3 = 1 - \frac{(1 - 2\nu_o)(1 + \nu_o + g\alpha_{11}^*)}{(1 + g\alpha_{33}^*)(1 - \nu_o + g\alpha_{11}^*) - 2\nu_o^2}$$

$$\text{with } g = \frac{32}{3} \frac{(1 - \nu_o^2)}{(2 - \nu_o)}$$



Berea sandstone data of
Lockner & Beeler (2003)



Confining Pressure: 30 MPa

Pore Pressure: 7.5 MPa

$E_o = 75$ GPa, $\nu_o = 0.15$

$C = 0.19$ GPa⁻¹ (from direct
storage measurement)

Skempton Tensor of an Anisotropically Cracked Rock:

Dependence on porosity and fracture anisotropy

- ❖ Using results of *Cheng* (1997) and *Sayers & Kachanov* (1995) the **storage coefficient** can be related to the crack density:

$$C = \frac{32(1 - \nu_o^2)}{3E_o(2 - \nu_o)} \alpha_{kk}^* + E_o \phi (\beta_f - \beta_s)$$

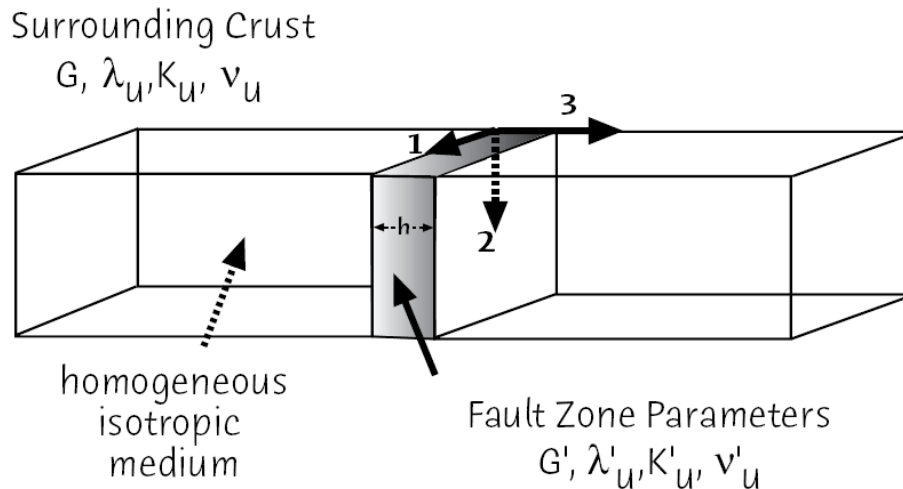
- ❖ Since typically $\beta_f \gg \beta_s$, we can neglect β_s and the principal components of the Skempton tensor are simply given by:

$$B_t = B_1 = B_2 = \frac{\alpha_{11}^*}{(\alpha_{kk}^* / 3) + \xi \phi}$$

$$B_z = B_3 = \frac{\alpha_{33}^*}{(\alpha_{kk}^* / 3) + \xi \phi}$$

with $\xi = \frac{E_o \beta_f (2 - \nu_o)}{32 (1 - \nu_o^2)}$ (**~2** for $\beta_f = 0.45 \text{ GPa}^{-1}$, $E_o = 75 \text{ GPa}$, $\nu_o = 0.15$)

Skempton Tensor of an Anisotropically Cracked Rock: Dependence on porosity and fracture anisotropy



$$B_t = B_1 = B_2 = \frac{\alpha_{11}^*}{(\alpha_{kk}^*/3) + \xi\phi}$$

$$B_z = B_3 = \frac{\alpha_{33}^*}{(\alpha_{kk}^*/3) + \xi\phi}$$

If $\alpha_{33}^* \gg \alpha_{11}^*$, then

$$B_3 = \frac{3}{1 + 3\xi\phi/\alpha_{33}^*} \gg B_1$$

For an array of cracks with aspect ratio c/a

$$B_3 = \frac{3}{1 + 4\pi\xi(c/a)} \text{ and } \Delta p = \frac{\Delta\sigma_n}{1 + 4\pi\xi(c/a)}$$

If $c \ll a$, then $B_3 \rightarrow 3$ and $\Delta p \sim \Delta\sigma_n$

Anisotropic Poroelasticity in Cracked and Granular Materials

- Under triaxial compression, appreciable stress-induced anisotropy develops in a rock, and as a result, its undrained mechanical response is anisotropic.
- Synthesizing the theoretical results of Cheng (1997) and Kachanov (1993), we derived explicit expressions for the Skempton pore pressure and Biot-Willis effective stress tensors in terms of the crack fabric tensors in a transversely isotropic material.
- Laboratory data on Berea sandstone show that the anisotropies associated with the Skempton and Biot-Willis coefficients increases with increasing differential stress and decreasing effective pressure, and accordingly the crack density is inferred to have similar anisotropic attributes.
- An analogous approach was also used to analyze laboratory data for Skempton tensor of Ottawa sand.